Entanglement generation in electron molecule inelastic scattering process



Motivation and Goal

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Collision processes are at the very core of many fields of physics, such as plasma physics, astrophysics, environmental science and radiation damage. On this work we discuss the generation and dynamics of entanglement during an inelastic collision (inter coulombic electron capture, ICEC) between an electron and a molecule, which may be dissociated after the collision. We computed the fully coupled electron-nuclear dynamics to describe the process.

Previous works addressed charge transfer and ionization processes using quantum information measures as witnesses of the dynamics. Here we compute quantum correlations developed during the inelastic process, quantified using several quantities such as quantum and shannon mutual information, shannon differential entropy and von Neumann entropy. We focus also on the correlations before, during and after the collision is over, and analyze the amount of the newly generated correlations that survive after the collision.

Methods

PECs calculation $\hat{H}_{\text{NeHe}^+}(R)\Phi_n(z_{e_1};R) = \epsilon_n(R)\Phi_n(z_{e_1};R)$ $\hat{W}_{\pm}^{(z_i)} = i\eta(z_i \mp z_{i\text{CAP}})^2\Theta(z_i \mp z_{i\text{CAP}}) \text{ CAPs}$ function $\psi_{\text{symm}}^{\text{init}} = \frac{1}{\sqrt{2}} [\phi_i(z_{e_1})\Phi_0(z_{e_2};R) + \phi_i(z_{e_2})\Phi_0(z_{e_1};R)]$ $\phi_i(z_{e_2})\Phi_0(z_{e_1};R)]$

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Symmetrized electronic spatial wave





$$\hat{H} = \sum_{i} \left[\frac{(\hat{\mathbf{p}}_{i})^{2}}{2m_{i}} + \hbar\omega_{i} \right] + \sum_{i,j \wedge i \neq j} V_{\text{eff}}^{(ij)}(|z_{ij}|)$$

$$V_{\text{eff}}^{(ij)}(|z_{ij}|) = \frac{\pi^{3/2}q_{i}q_{j}}{2\left(\sqrt{2}l_{c}\right)} \exp\left[\left(\frac{|z_{ij}|}{\sqrt{2}l_{c}} \right)^{2} + \left(\frac{\sqrt{2}l_{c}}{2l_{\alpha_{ij}}} \right)^{2} \right] \operatorname{erfc}\left(\frac{|z_{ij}|}{\sqrt{2}l_{c}} + \frac{\sqrt{2}l_{c}}{2l_{\alpha_{ij}}} \right)$$

von Neumann entropy

 $S_i^{\rm vN}(t) = -\sum_{\alpha=1}^N \lambda_\alpha^{(i)}(t) \log_2\left(\lambda_\alpha^{(i)}(t)\right)$

where i, j, k can be $\{e_1, e_2, c_{12}\}$

but different from each other

Negative values of the CEs are a sufficient (not necessary)

condition for the presence of entanglement [45]. The negative

Entropies

Shannon differential entropy $S_{z_{i}}^{\text{Sh}}(t) = \int \rho(z_{i}, t) \log_{2}[\rho(z_{i}, t)] dz_{i} \qquad S_{z_{i}}^{\text{Sh}}(t) = \int \int \rho(z_{i}, z_{j}, t) \log_{2}[\rho(z_{i}, z_{j}, t)] dz_{i} dz_{j}$ $Von \text{Neumann conditional entropy} \qquad S_{\alpha_{1}|\alpha_{2}}^{\text{vN}}(t) = S_{\alpha_{1},\alpha_{2}}^{\text{vN}}(t) - S_{\alpha_{2}}^{\text{vN}}(t) \qquad \text{cond} \\ S_{\alpha_{1},\alpha_{2}|\alpha_{3}}^{\text{vN}}(t) = S_{\alpha_{1},\alpha_{2},\alpha_{3}}^{\text{vN}}(t) - S_{\alpha_{3}}^{\text{vN}}(t) \qquad \text{ment} \\ S_{e|N}^{\text{vN}}(t) = -S_{e|e}^{\text{vN}}(t) = S_{e}^{\text{vN}}(t) - S_{N}^{\text{vN}}(t) \qquad \text{ment} \\ \end{array}$

 $S_{N|e}^{\rm vN}(t) = 0$

 $S_{e,e|N}^{\mathrm{vN}}(t) = -S_N^{\mathrm{vN}}(t)$

 $a_{\alpha_{3}}(t) - S_{\alpha_{3}}^{vN}(t)$ $= S_{e}^{vN}(t) - S_{N}^{vN}(t)$ value obtained in Eq. (33) indicate the presence of entanglement between electrons (both taken as one part) and the nuclei coordinate, while Eq. (34) indicate the presence of entanglement between one electron and the nuclei-electron compound. The sign of Eq. (31) is not defined *a priori*; however, we will show that for the system and processes considered here, $S_{e}^{vN}(t) > S_{N}^{vN}(t)$ for all *t* (see Figs. 4–6), and thus the entanglement between electrons is confirmed [$S_{e|e}^{vN}(t) < 0$]. **Mutual informations**

$\begin{array}{ll} S_{e,N|e}^{\mathrm{vN}}(t) = -S_{e}^{\mathrm{vN}}(t) & \mathbf{Mutual informations} \\ \text{The QMI measures correlations (quantum and classical)} \\ \text{between subsystems. It is non-negative (due to sub-} I_{e,e}^{\mathrm{vN}}(t) : I_{e,e|N}^{\mathrm{vN}}(t) = 2S_{e|N}^{\mathrm{vN}}(t) - S_{e,e|N}^{\mathrm{vN}}(t) \\ \text{additivity of VNE}). QMI exceeds the bound of classical } = 2S_{e}^{\mathrm{vN}}(t) - S_{N|e}^{\mathrm{vN}}(t) \\ \text{MI, because quantum systems can be supercorrelated.} & I_{e,N|e}^{\mathrm{vN}}(t) : I_{e,N|e}^{\mathrm{vN}}(t) = S_{e|e}^{\mathrm{vN}}(t) - S_{e,N|e}^{\mathrm{vN}}(t) \\ \text{The SMIs represent the degree of correlations between} & = S_{N}^{\mathrm{vN}}(t) \\ \end{array}$

